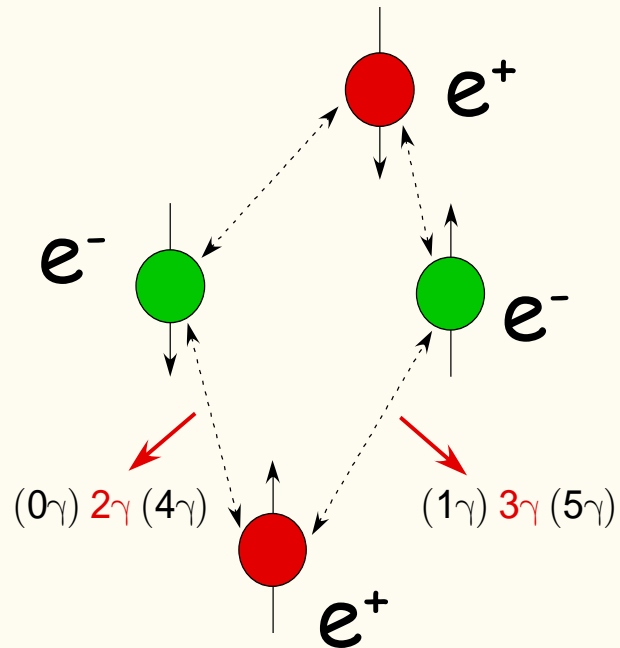


Dipole excitation of dipositronium



Few-Body Problems in Physics

Bonn, August 31, 2009

Andrzej Czarnecki  University of Alberta

Outline

Positronium, its ion, and the molecule

1951

1981

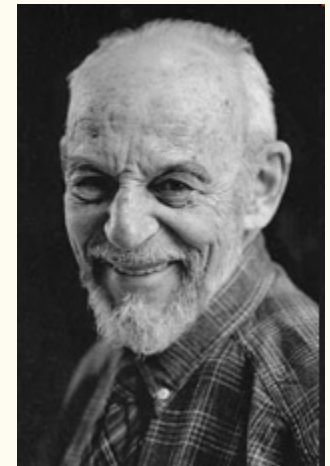
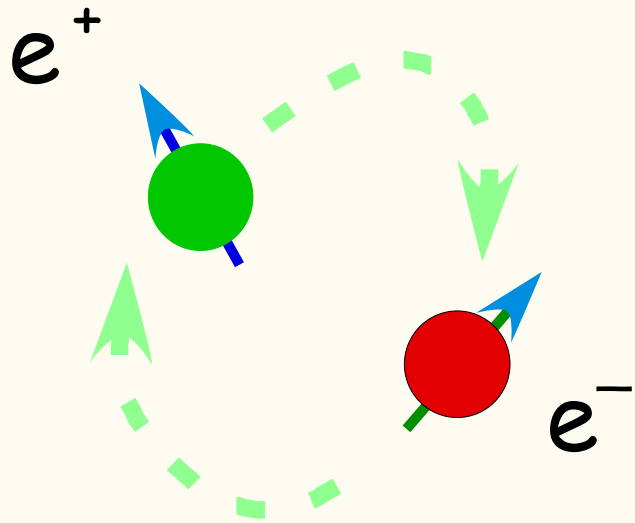
2007

Hyperfine splitting in Ps: new measurement

Ps⁻ ion: theoretical prediction for the decay rate. Comparison with experiment.

Molecule: spectrum. Dipole transition vs. annihilation.

Positronium



Martin Deutsch
1917 - 2002

Very similar to hydrogen, except

- no hadronic nucleus
- annihilation
- reduced mass reduced

$$m_e \rightarrow \frac{m_e}{2}$$

Two spin states:

singlet (para-Ps; short-lived, 0.1 ns)

triplet (ortho-Ps; long-lived, ~150 ns)

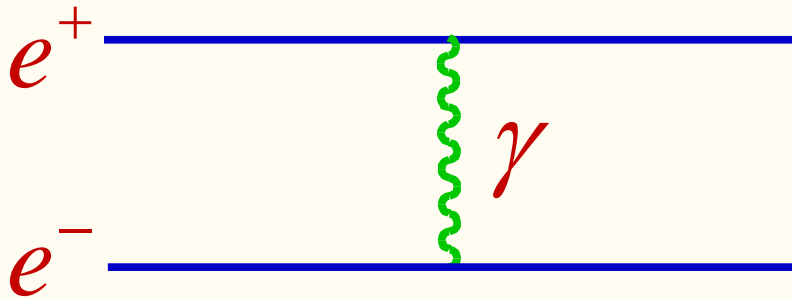
Binding: 0.25 a.u.

All properties can be described by QED, using one parameter: $\alpha = \frac{1}{137.036}$

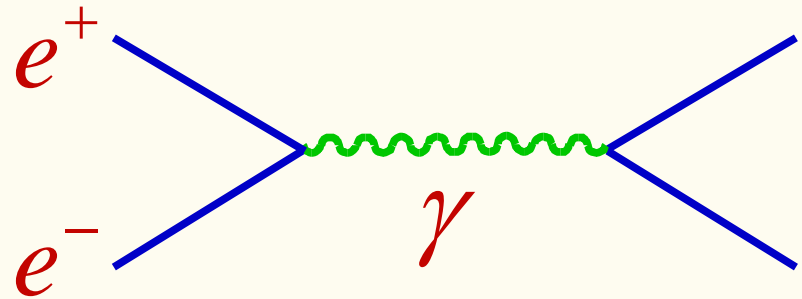
Positronium spectrum: discrepancy with QED

Tree-level QED prediction for the hyperfine splitting (HFS)

Recoil effect



Annihilation effect

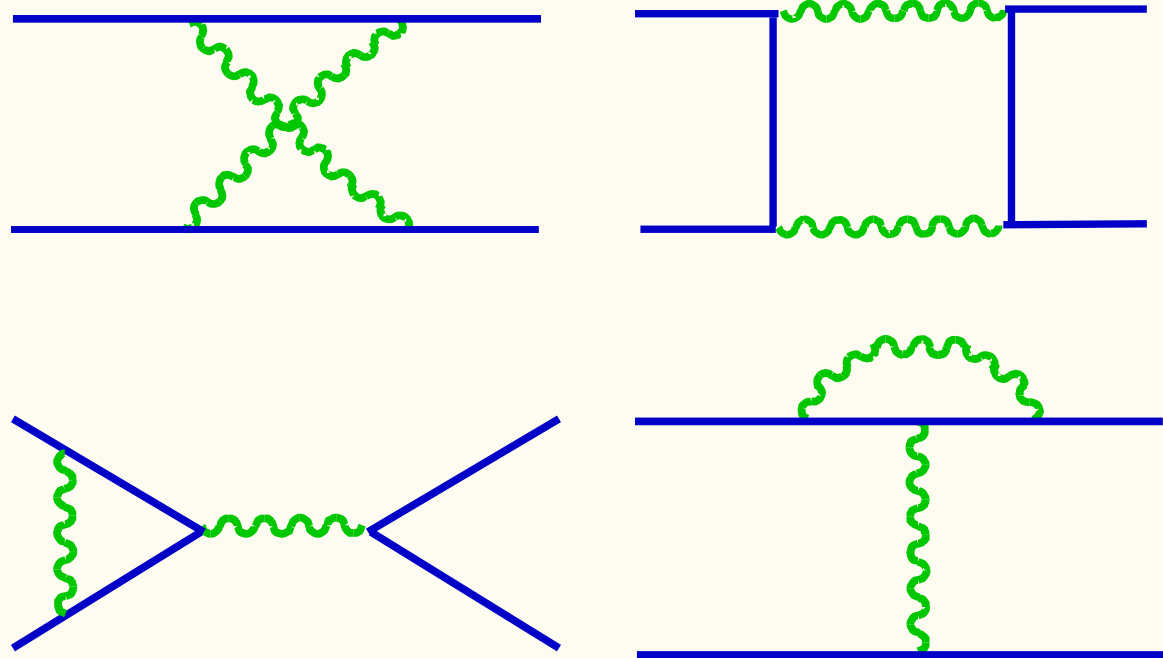


$$\gamma_{\mu} \otimes \gamma^{\mu} \rightarrow 1 \otimes 1 + \sigma \otimes \sigma$$

$$\Delta \nu_{\text{HFS}} = \frac{7}{12} m_e \alpha^4 \simeq 204 \text{ GHz}$$

Quantum corrections to the HFS: one-loop

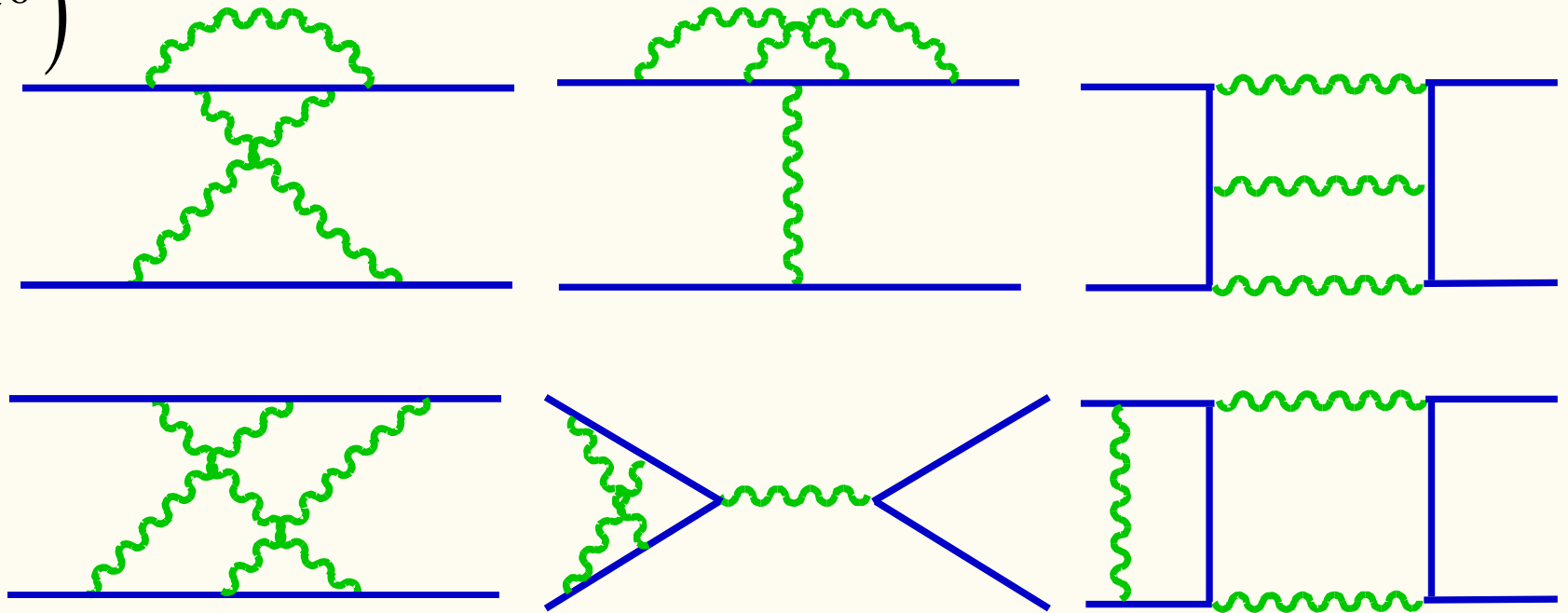
$O(\alpha^5)$



$$-\frac{m_e \alpha^5}{\pi} \left(\frac{8}{9} + \frac{\ln 2}{2} \right) \simeq -1005.5 \text{ MHz} \rightarrow -0.5\%$$

Quantum corrections to the HFS: two-loop

$O(\alpha^6)$

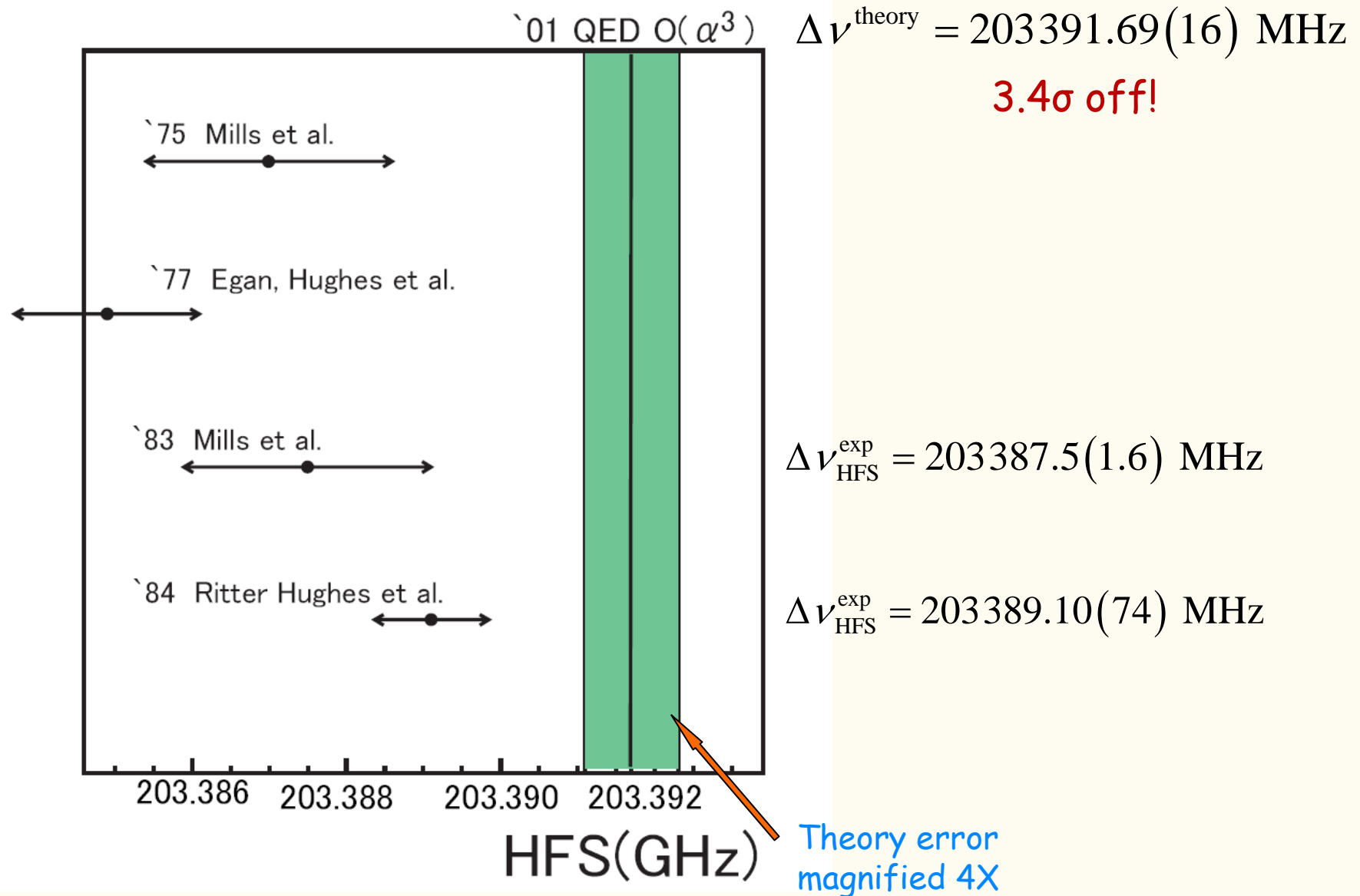


Recoil

$$\frac{m_e \alpha^6}{\pi^2} \left[\frac{1367}{648} - \frac{5197}{3456} \pi^2 + \left(\frac{1}{2} + \frac{221}{144} \pi^2 \right) \ln 2 - \frac{53}{32} \zeta(3) + \frac{5}{24} \pi^2 \ln \frac{1}{\alpha} \right]$$

$\simeq 11.8 \text{ MHz} \rightarrow 0.006\%$ (Experimental error $\simeq 0.7 \text{ MHz}$)

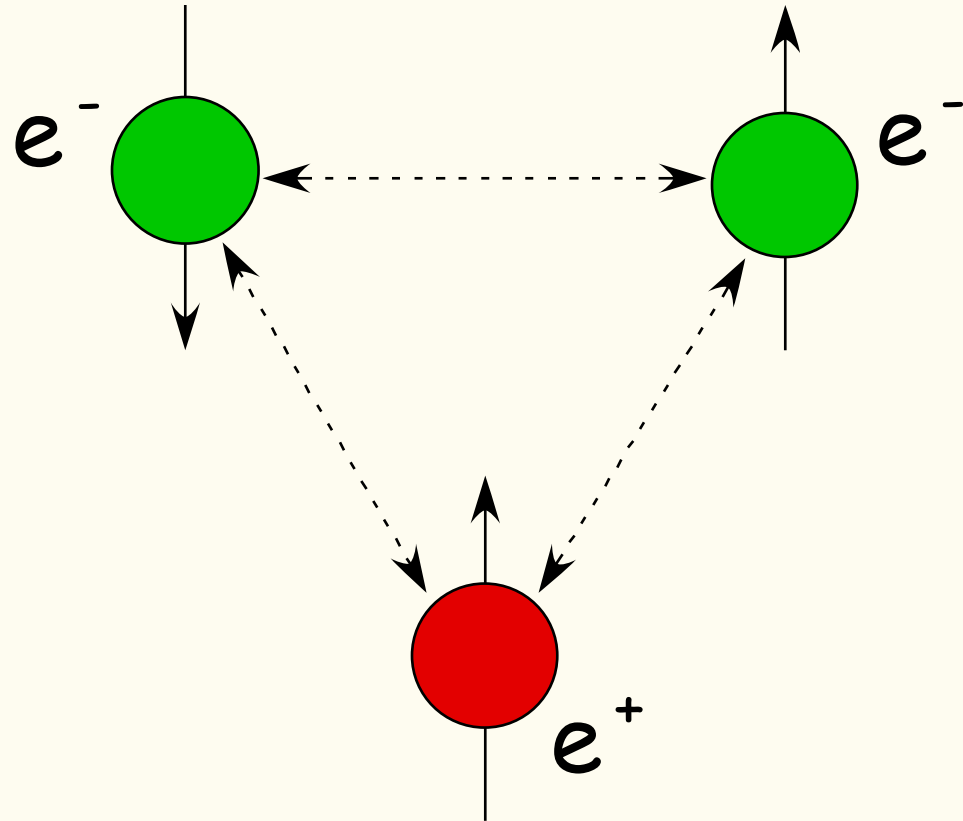
HFS theory vs. measurements



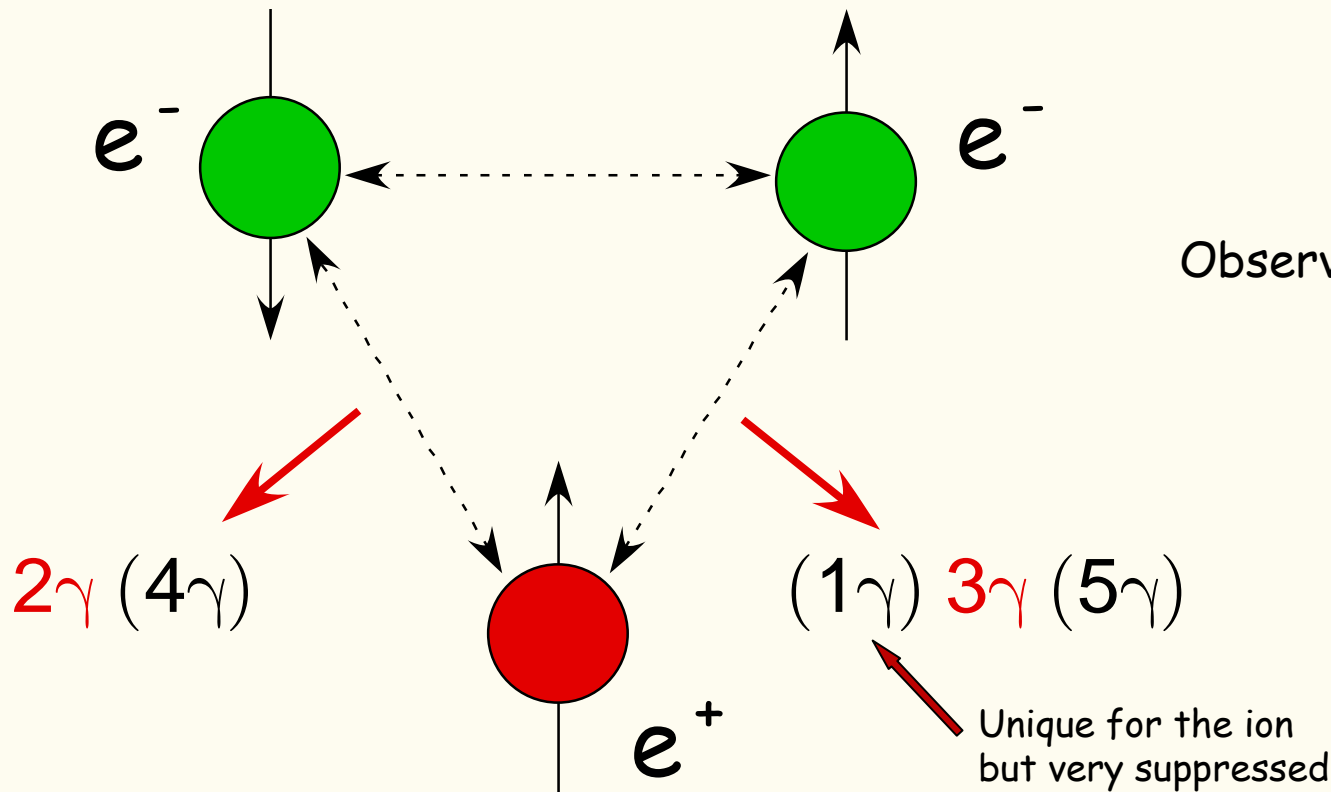
Positronium ion: predicted in 1946



John Wheeler
(1911-2008)



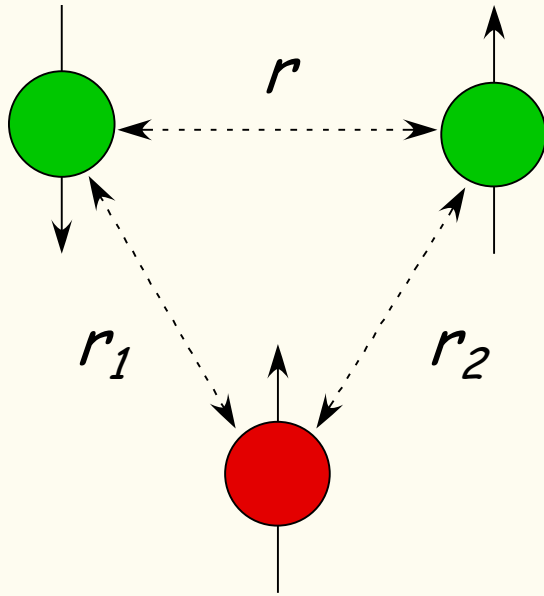
Positronium ion: a new test of bound-state QED



New positronium-ion source
at the FRM II reactor in Munich:
measurements of branching ratios.



Theory of the Ps ion: the wavefunction



$$\psi(\vec{r}_1, \vec{r}_2) = \frac{\uparrow\downarrow - \downarrow\uparrow}{\sqrt{2}} \phi(r_1, r_2, r)$$

The wave function is not known analytically, but can be found using the variational method.

The ion resembles a positronium "shell" and a loosely-bound electron.

Now we can easily estimate the dominant decay channels.

Recent measurement of the ion decay rate

63401 (2006)

PHYSICAL REVIEW LETTERS

week
17 FEBRU

Measurement of the Decay Rate of the Negative Ion of Positronium (Ps^-)

Frank Fleischer,^{1,*} Kai Degreif,¹ Gerald Gwinner,^{1,†} Michael Lestinsky,¹
Vitaly Liechtenstein,² Florian Plenge,¹ and Dirk Schwalm¹

¹*Max-Planck-Institut für Kernphysik, Heidelberg, Germany*

²*Kurchatov Institute, Moscow, Russia*

(Received 28 November 2005; published 13 February 2006)

A new determination of the decay rate of the negative ion of positronium (Ps^-), using a beam-foil method and a stripping-based detection technique, is reported. The measured result of $\Gamma = 2.089(15) \text{ ns}^{-1}$ is a factor of 6 more precise than the previous experimental value of $\Gamma = 2.09(9) \text{ ns}^{-1}$, and is in excellent agreement with the theoretical value of $\Gamma = 2.086(6) \text{ ns}^{-1}$.

$$\Gamma_{\text{exp}}(\text{Ps}^-) = 2089(15) \mu\text{s}^{-1}$$

$7 \cdot 10^{-3}$
(factor 4-5
improvement
expected)

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63401 (2006)

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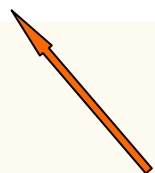
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$$\Gamma_{\text{exp}}(\text{Ps}^-) = 2089(15) \mu\text{s}^{-1}$$

$$7 \cdot 10^{-3}$$

(factor 4-5
improvement
expected)



We wanted to
improve this

Positronium ion decay: refinements

Corrections $O(\alpha)$

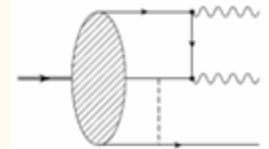
single hard photon loops

Corrections $O(\alpha^2)$; challenge: divergences $\rightarrow \ln \alpha$

soft

$O(k^2)$ corrections to the amplitude M

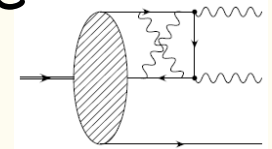
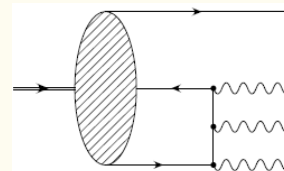
Breit hamiltonian \rightarrow correction to $\psi(r=0)$



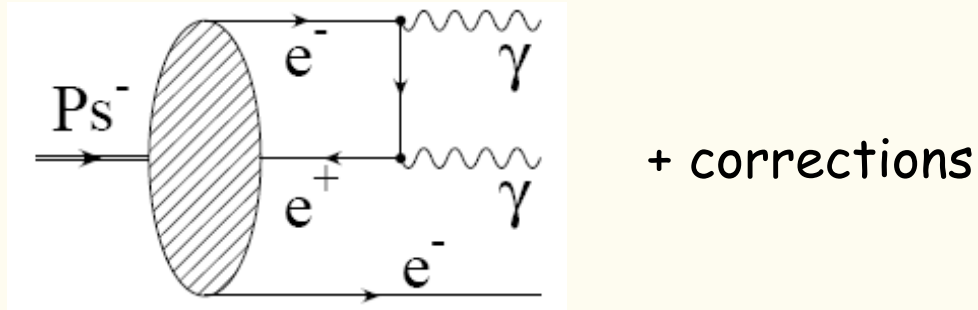
hard

Short-distance two-loop photon exchange

Real photon radiation



Decay rate prediction

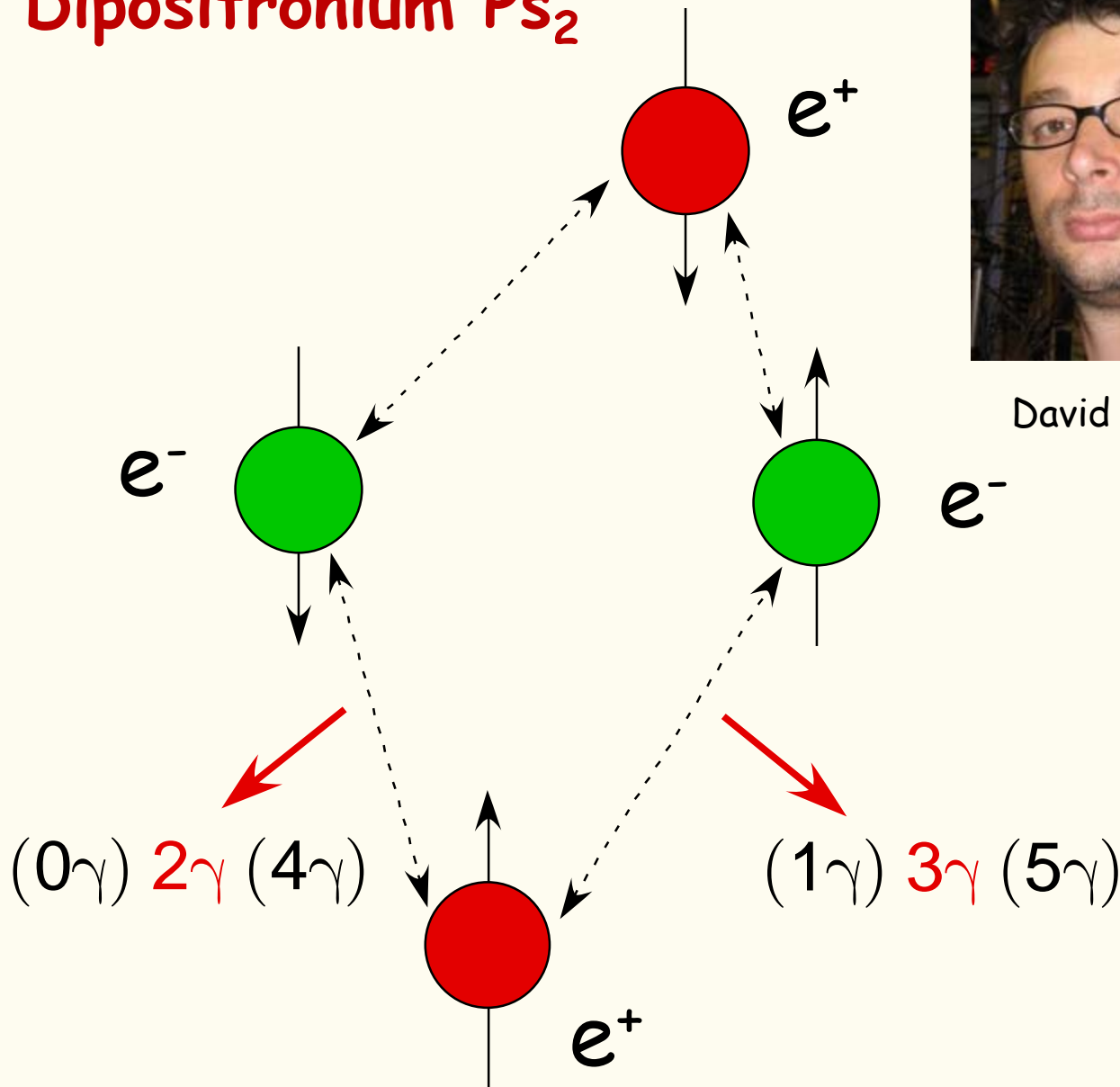


$$\Gamma(\text{Ps}^-) = 2\pi \frac{\alpha^5 m_e c^2}{\hbar} (1 + C) \langle \delta^3(r_{12}) \rangle$$

$$\Gamma(\text{Ps}^-) = 2.087963(12) \text{ ns}^{-1}$$

with M. Puchalski and S. Karshenboim
PRL **99**, 203401 (2007)

Dipositronium Ps_2

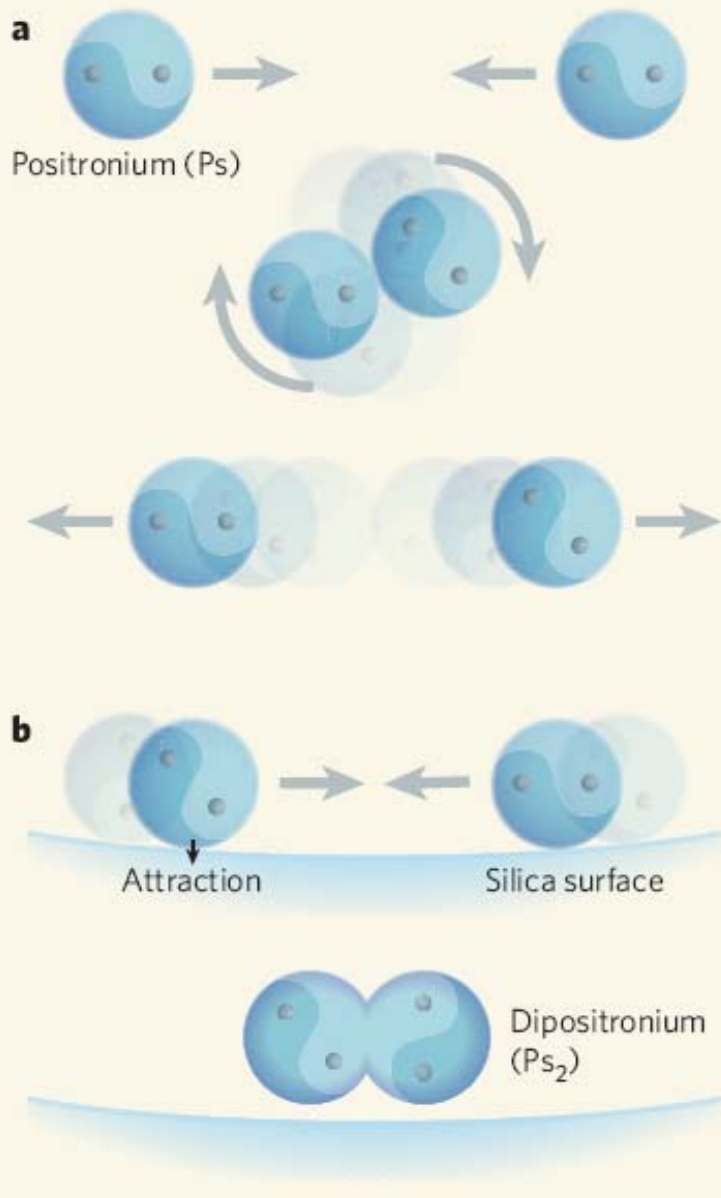


David Cassidy



Allen Mills

Discovery of dipositronium 2007



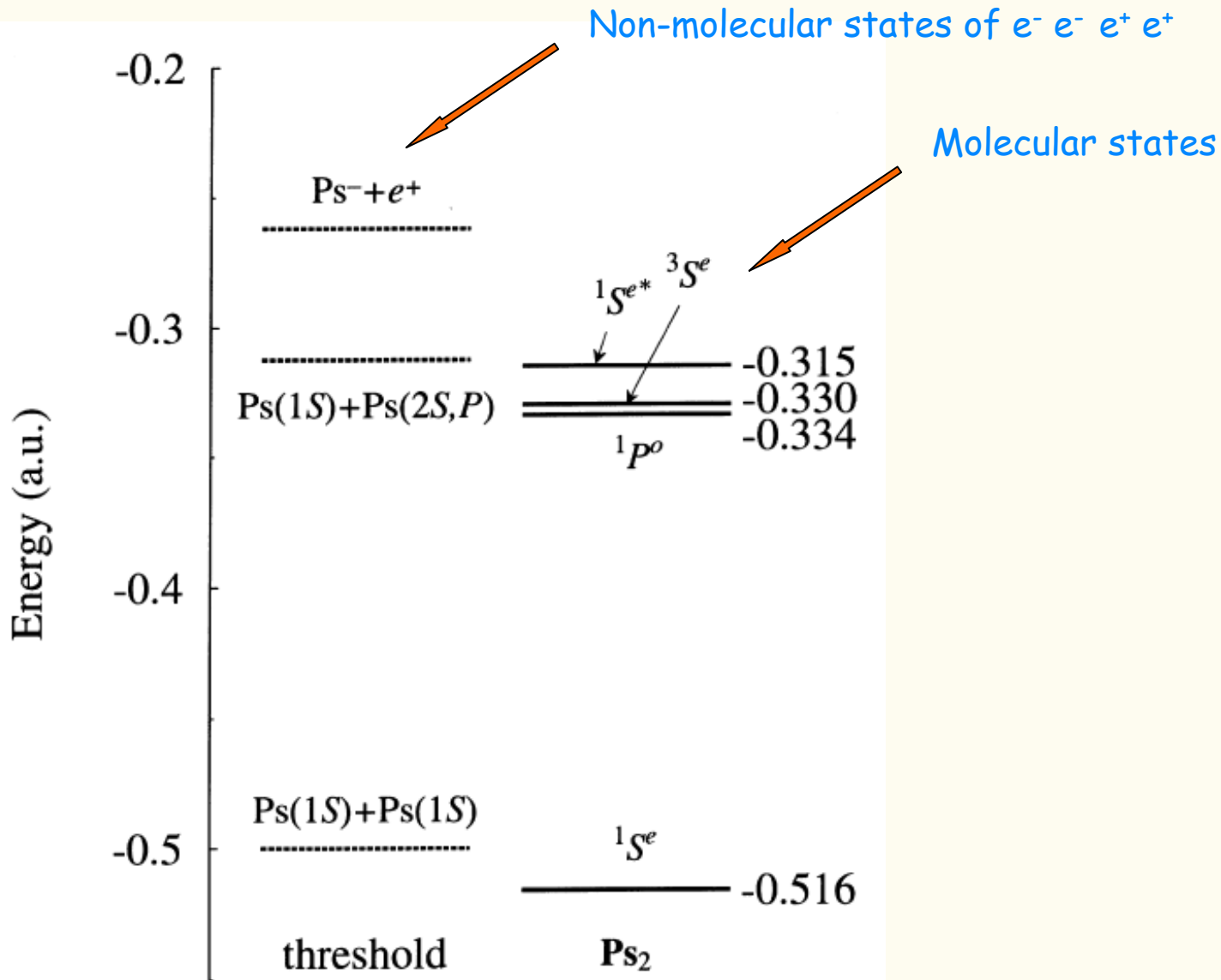
Molecule formation kills long-lived positronia.

At higher temperature, fewer atoms on the surface, fewer molecules formed.

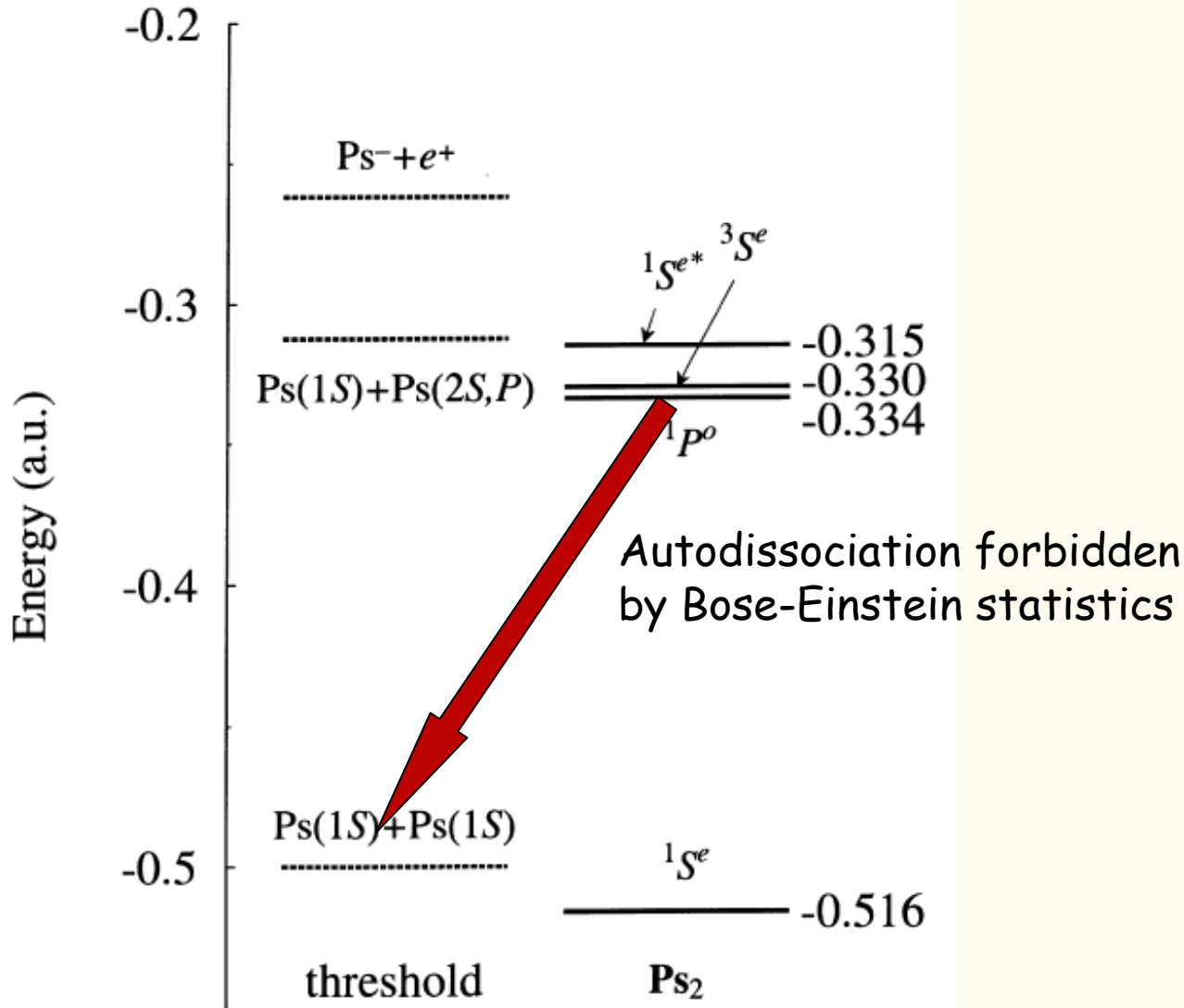
Indeed: at high-T, more long-lived positronia observed.

Cassidy & Mills, Nature 2007

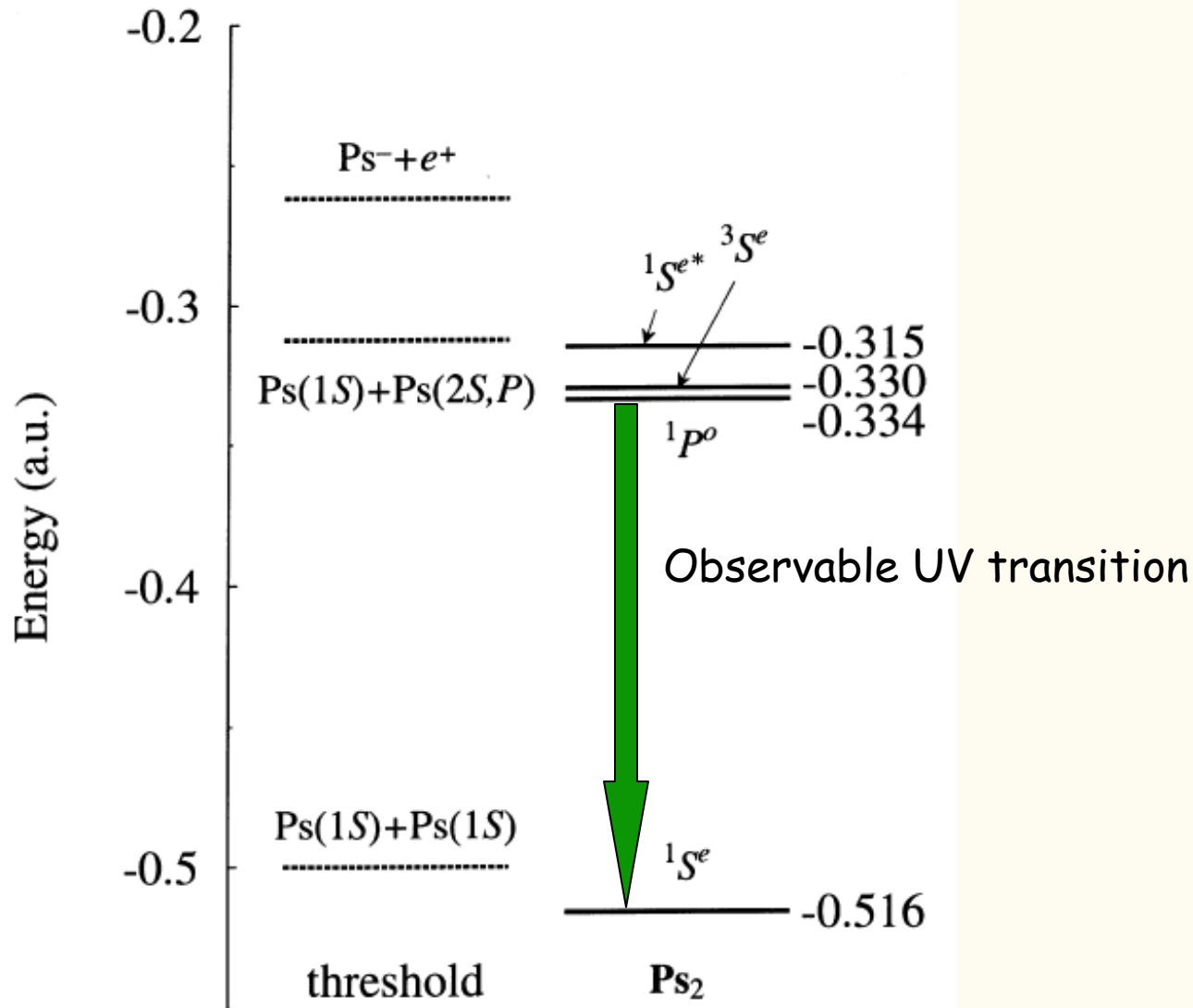
Spectrum of the molecule Ps_2



A direct signal of the molecule: transition line.



A direct signal of the molecule: transition line.



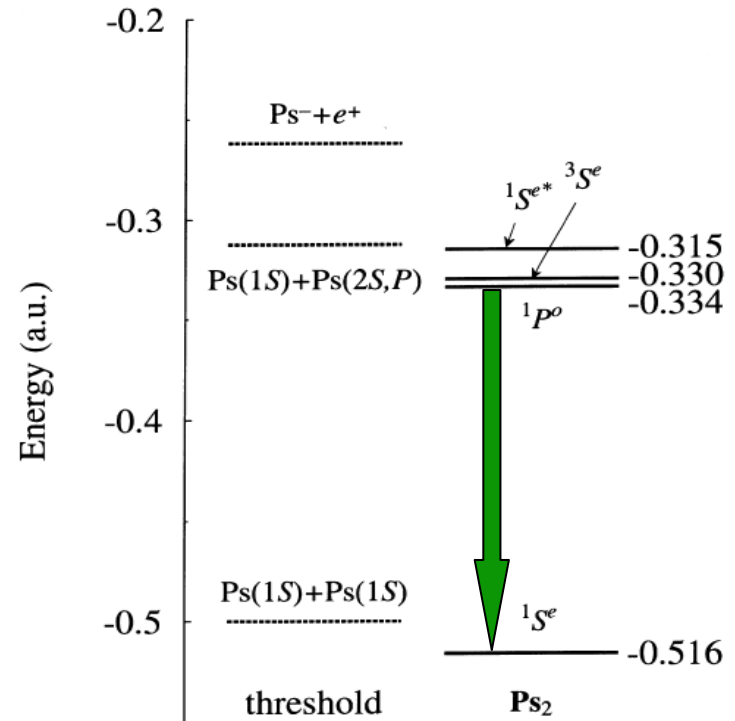
Questions about this transition:

What is its accurate energy?

$$\Delta E = E_P - E_S = 0.1815867(8) \text{ a.u.} \\ \simeq 4.9 \text{ eV}$$

Similar to atomic positronium,
but softer (dielectric effect?):

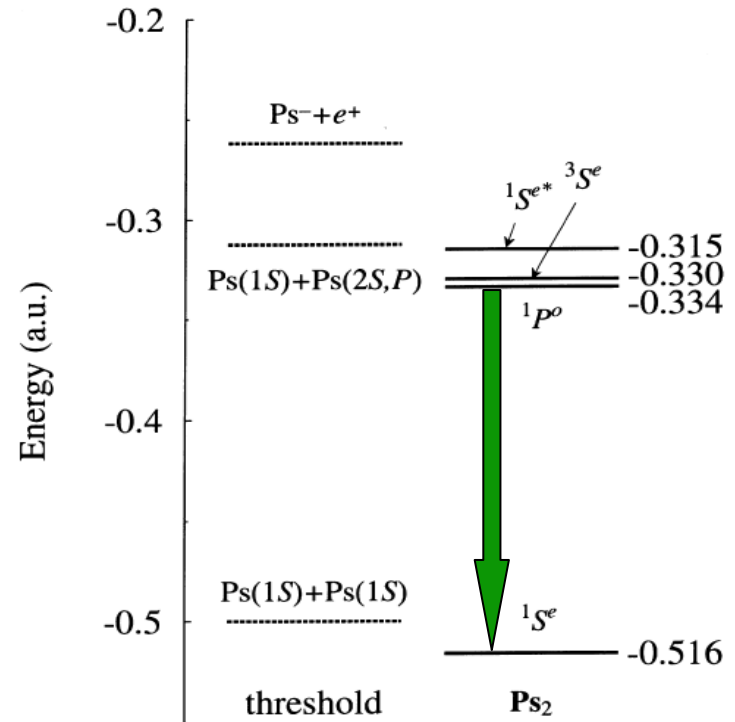
$$E_P - E_S = \frac{3}{4} \times \frac{1}{4} \text{ a.u.} = 0.1875 \text{ a.u.}$$



Questions about this transition:

How often does
radiative transition appear
(before annihilation)?

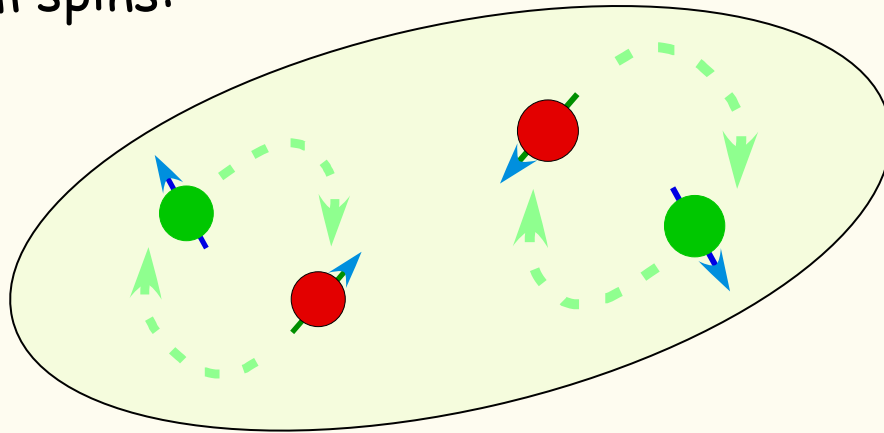
$$\text{BR}(P \rightarrow S) = \frac{\Gamma_{\text{dip}}(P \rightarrow S)}{\Gamma_{\text{annih}}(P) + \Gamma_{\text{dip}}(P \rightarrow S)} = 0.191(2)$$



Competition: dipole transition vs. annihilation

S-state annihilates quite rapidly.

Assume it consists of two weakly-interacting Ps atoms, with random spins:



There is a para-positronium pair with probability

$$2_{\text{pairs}} \cdot \frac{1}{4_{\text{spin states}}} = \frac{1}{2}$$

The decay rate: $\Gamma_s \approx \frac{1}{2_{\text{Ps probability}}} \cdot \frac{1}{\tau_{\text{pPs}}} = \frac{1}{0.25 \text{ ns}}$

P-state: half of this rate: $\Gamma_{\text{annih}} = \frac{1}{0.5 \text{ ns}}$

Competition: dipole transition vs. annihilation

In an isolated P-excited Ps atom lives ~ 3.2 ns;
In Ps_2 the decay is about twice faster
(and the same as in atomic hydrogen)

$$d_{\text{atom}} = \langle S | \vec{d} | P \rangle$$

$$d_{\text{molecule}} = \left\langle SS \left| \vec{d} \right| \frac{PS+SP}{\sqrt{2}} \right\rangle = \sqrt{2} d_{\text{atom}}$$

$$\Gamma_{\text{dip}}(P \rightarrow S) = (\sqrt{2})^2 \Gamma_{\text{dip}}(\text{atomic Ps}) = \frac{1}{1.6 \text{ ns}}$$

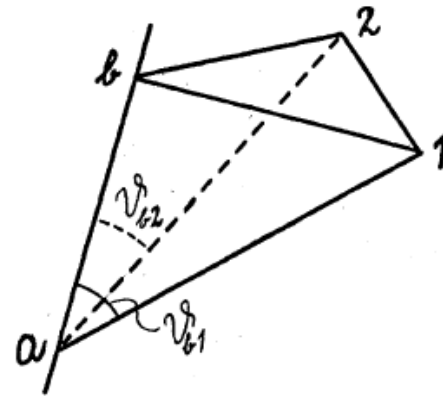
Branching ratio

$$BR(P \rightarrow S) = \frac{\Gamma_{\text{dip}}(P \rightarrow S)}{\Gamma_{\text{annih}}(P) + \Gamma_{\text{dip}}(P \rightarrow S)} \sim \frac{\frac{1}{1.6}}{\frac{1}{1.6} + \frac{1}{0.5}} \sim 19\%$$

Variational determination of the Ps_2 wave fnc.

Gaussian basis

$$\exp\left(-\sum_{i=1}^6 a_i r_i^2\right)$$



Hylleraas & Ore, 1947

Coordinate system for the positronium molecule

Relatively small basis ~ 2000

QR decomposition of the eigenvalue equation

Optimization of individual basis elements with Powell's method

\rightarrow nonrelativistic energies ~ 1 ppb

(test: Lithium)

Relativistic corrections dominated by annihilation: repulsive,
more effective in the ground state \rightarrow decreases the SP interval.

Summary

Methods of bound-state QED developed largely with positronium.

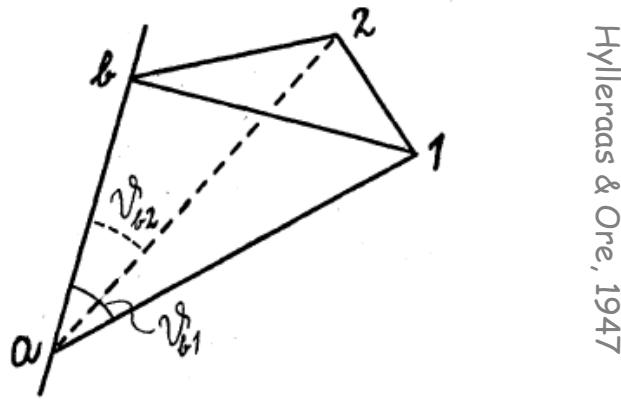
New states of positronium-like systems, ions and molecules: opportunity to test three- and four-body QED.

The Coulomb wave functions not known analytically for $e^+ e^- e^-$ and $e^+ e^- e^+ e^-$.

Variational methods very accurate. $O(\alpha^2)$ corrections computed for the ion decay rate and the Ps_2 P-S interval. Both are being measured.

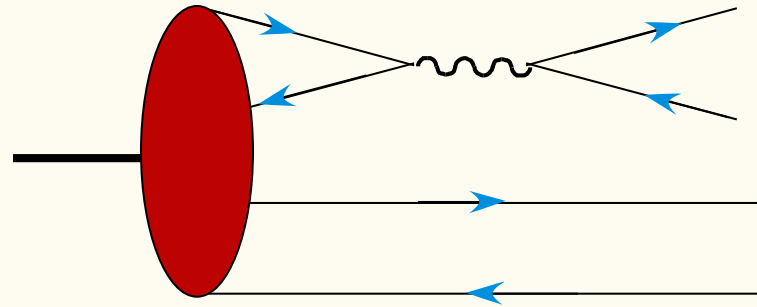
Energy levels: ground state and P-excitation

Wave function determined variationally, using Coulomb potential; M. Puchalski



Coordinate system for the positronium molecule

Relativistic corrections: perturbations.
Annihilation dominates.



Interval P-S determined with 5×10^{-6} accuracy (slightly smaller than in Ps, "dielectric effect").

Breakdown of corrections to the Ps ion width

Correction	Value
$\alpha A^{3\gamma}$	0.002 693 245
$\alpha A^{2\gamma}$	-0.005 882 770
$-2\alpha^2 \ln\alpha$	0.000 524 019
$\alpha^2 B^{4\gamma}$	0.000 001 480
$\alpha^2 B^{3\gamma}$	-0.000 064 352
$\alpha^2 B_{\text{squared}}$	0.000 008 652
$\alpha^2 B_{\text{hard}}^{\text{fin}}$	-0.000 218 3(34)
$\alpha^2 B_{\text{aa}}$	0.000 017 750
$\alpha^2 B_{\text{H1}}$	0.000 078 366
$\alpha^2 B_{\text{H2}}$	0.000 541 484
$3\alpha^3 \ln^2\alpha / (2\pi)$	-0.000 004 491
$2.5(2.5)\alpha^3 \ln\alpha$	-0.000 004 8(48)
Total C	-0.002 309 7(59)

$$\Gamma(\text{Ps}^-) = 2\pi \frac{\alpha^5 m_e c^2}{\hbar} (1 + C) \langle \delta^3(r_{12}) \rangle$$

Breakdown of corrections to the Ps_2 interval

$$H_{\text{rel}} = H_{\text{MV}} + H_{\text{D}} + H_{\text{OO}} + H_{\text{SS}} + H_{\text{A}},$$

$$H_{\text{MV}} = -\frac{1}{8} \sum_a \tilde{p}_a^4,$$

$$H_{\text{D}} = -\pi \sum_{a<b} z_{ab} \delta^3(r_{ab}),$$

$$H_{\text{OO}} = -\frac{1}{2} \sum_{a<b} z_{ab} p_a^i \left(\frac{\delta^{ij}}{r_{ab}} + \frac{r_{ab}^i r_{ab}^j}{r_{ab}^3} \right) p_b^j,$$

$$H_{\text{SS}} = -\frac{2\pi}{3} \sum_{a<b} z_{ab} \vec{\sigma}_a \cdot \vec{\sigma}_b \delta^3(r_{ab}),$$

$$H_{\text{A}} = \frac{\pi}{2} \sum_{a<b, ab \neq 12, 34} (3 + \vec{\sigma}_a \cdot \vec{\sigma}_b) \delta^3(r_{ab}).$$

Source	Ground state	P state
H_{C}	-0.516 003 790 415(88)	-0.334 408 317 34(81)
$\alpha^2 H_{\text{MV}}$	-0.000 009 152	-0.000 004 780(1)
$\alpha^2 H_{\text{OO}}$	-0.000 013 470	-0.000 007 736
$\alpha^2 H_{\text{D}}$	0.000 014 592	0.000 007 458
$\alpha^2 H_{\text{SS}}$	0.000 000 419	0.000 000 097
$\alpha^2 H_{\text{A}}$	0.000 022 202	0.000 011 259
$\alpha^2 H_{\text{rel}}$	0.000 014 591	0.000 006 298(1)
$\alpha^3 \ln \alpha H_{\text{log}}$	0.000 001 01(50)	0.000 000 51(25)
Total	-0.515 988 2(5)	-0.334 401 5(3)
[12]	-0.515 989 199 656	

$$\Delta E \equiv E_P - E_S = 0.181 586 7(8) \text{ a.u.}$$

Ratio of three- to two-photon annihilation

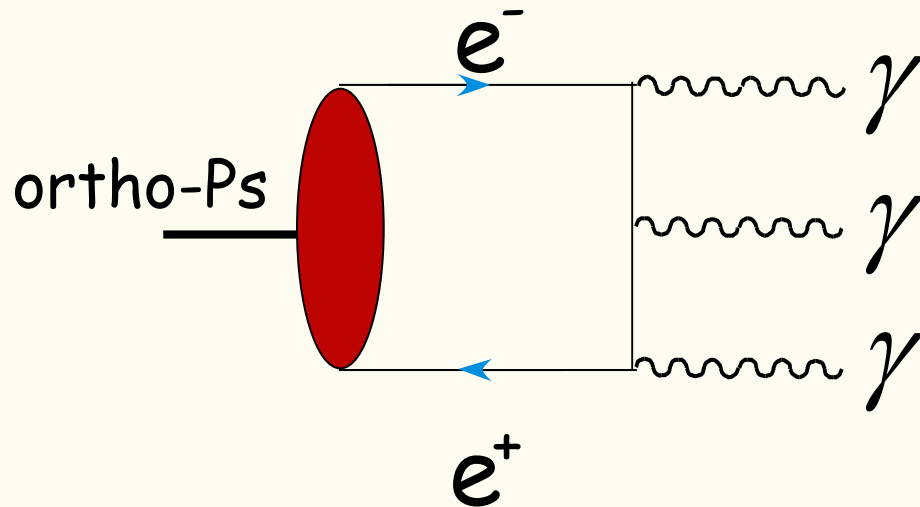
$$\text{BR}(\text{Ps}^- \rightarrow \gamma\gamma\gamma) \equiv \frac{\Gamma(\text{Ps}^- \rightarrow \gamma\gamma\gamma)}{\Gamma(\text{Ps}^-)}$$

$$= \alpha \left[A^{3\gamma} + \alpha(B^{3\gamma} - AA^{3\gamma}) - \frac{7}{3} A^{3\gamma} \alpha^2 \ln \frac{1}{\alpha} + \dots \right]$$

$$= 0.002\,635\,8(8).$$

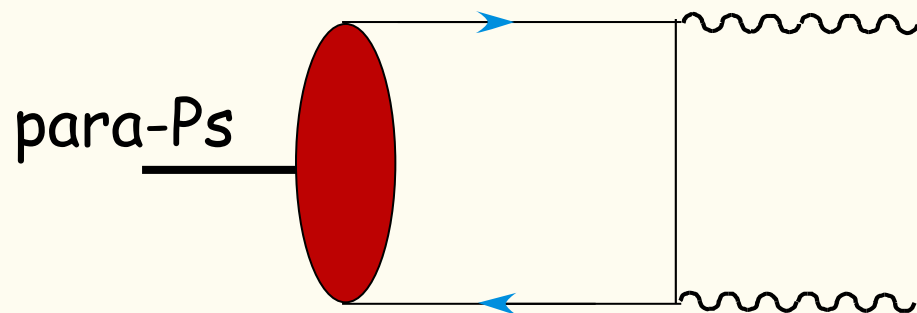
Back to positronium...

... and its decays



$$\Gamma_{\text{exp}}(\text{o-Ps}) = 7.0401(7) \mu\text{s}^{-1}$$

$$1 \cdot 10^{-4}$$



$$\Gamma_{\text{exp}}(\text{p-Ps}) = 7.9909(17) \text{ns}^{-1}$$

$$2 \cdot 10^{-4}$$

Theory of positronium decay

Corrections $O(\alpha)$

single hard photon loops

Corrections $O(\alpha^2)$

Para: AC, Melnikov, Yelkhovsky
Ortho: Adkins, Fell, Sapirstein

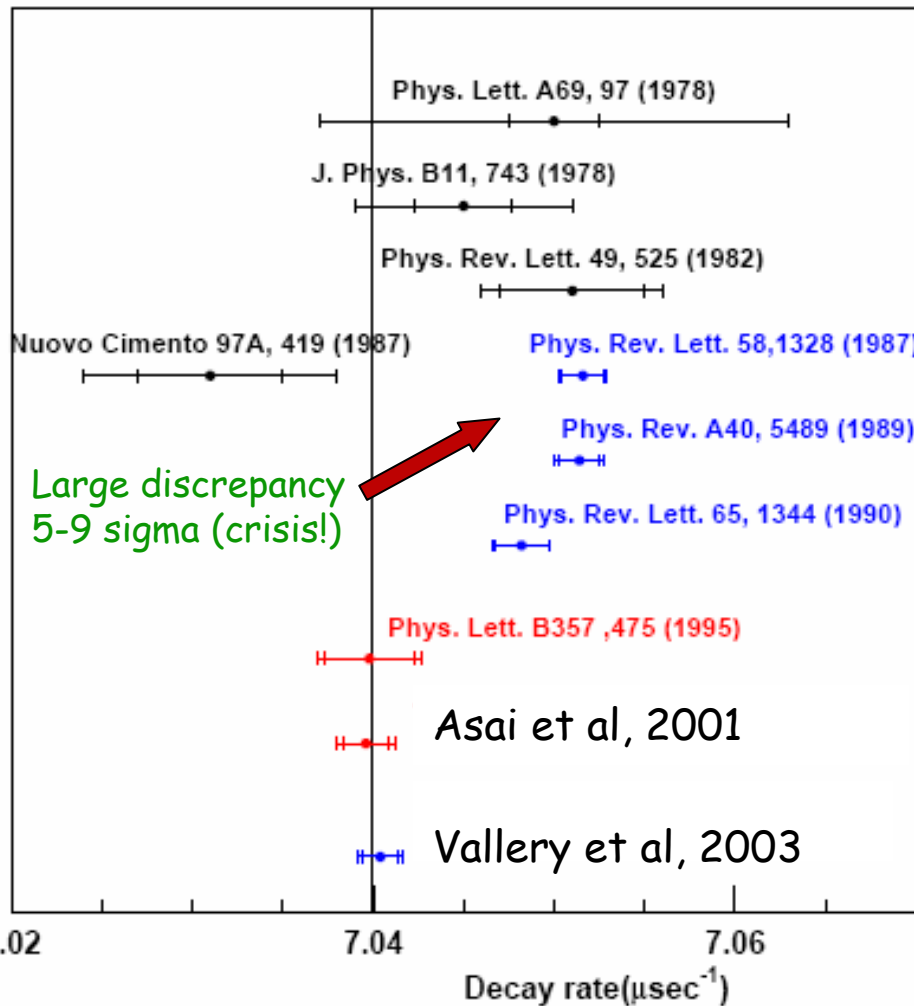
soft $\left\{ \begin{array}{l} O(k^2) \text{ corrections to decay amplitude} \\ \text{Breit hamiltonian} \rightarrow \text{correction to } \psi(r=0) \end{array} \right\}$

hard $\left\{ \begin{array}{l} \text{Short-distance two-loop photon exchange} \end{array} \right\}$

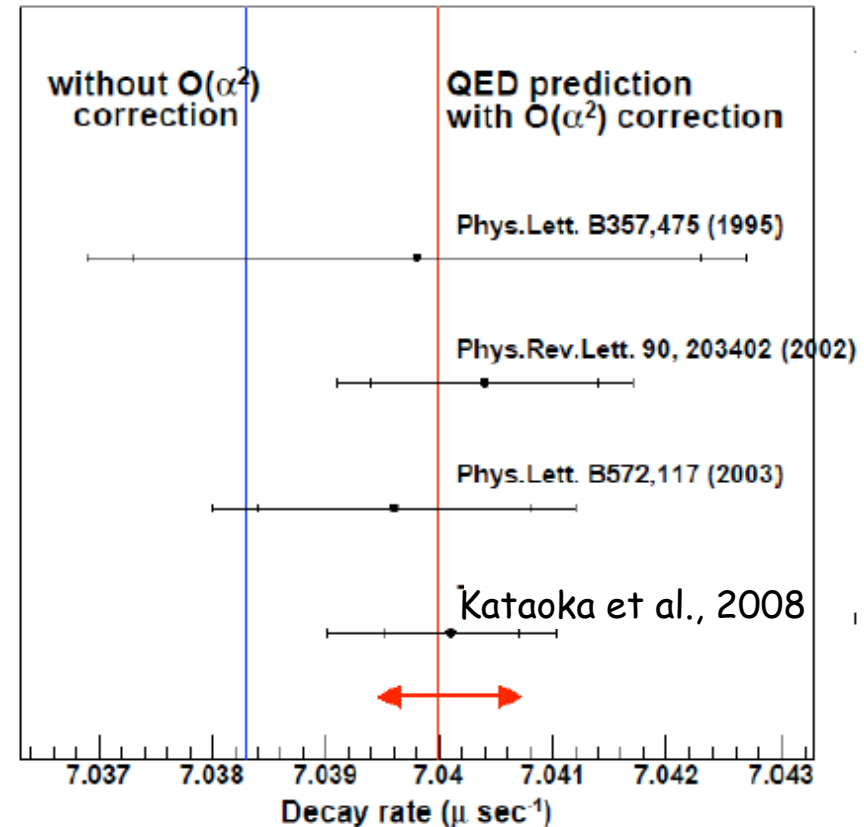
finite together, but give $\ln \frac{m}{m\alpha} = \ln \frac{1}{\alpha}$

History of oP s lifetime measurements

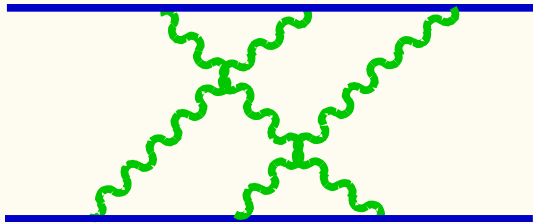
Older:



Latest ones, enlarged scale:



Recoil effect and $\log(\alpha)$



Hard photons, $k \sim m_e$, generate a $\delta^3(r)$ potential. Its coefficient can be found from scattering of **free** electrons: it is IR divergent.

Conversely, soft photons, through the iterated Breit interaction, give a UV divergence.

The trick is to find a common regularization scheme for both.

Pineda and Soto suggested (1998) dimensional regularization,

$$\frac{(\alpha m)^{d-3}}{d-3} - \frac{m^{d-3}}{d-3} \xrightarrow{d \rightarrow 3} \ln \alpha$$